As many fish scales serve as efficient natural dermal armors that protect fish from predators without compromising their flexibility, mimicking their design may lead to improved lightweight synthetic armors. The scales of the common carp have evolved with a mineralized outer layer to resist penetration and a tougher lower layer with a twisted arrangement of mineralized collagen fibrils, which absorbs excessive deformation through a sequence of multiple plasticity mechanisms.
Structure and Mechanical Adaptability of a Modern Elasmoid Fish Scale from the Common Carp

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SUMMARY
The carp (Cyprinus carpio) has typical elasmoid scales commonly found on teleosts. They provide protection while retaining flexibility and maneuverability of the fish. The exterior surface of the scale consists of an ultrathin discontinuous mineral layer on top of mineralized woven collagen fibrils. The underlying foundation is composed of two collagenous components. The major one consists of a single-twisted “Bouligand” structure with a twisting angle of 36°. A secondary “sheet-like” structure, formed by thinner collagen fibrils oriented along the thickness direction, acts to increase the integrity of the scale. Here, we identify the deformation and failure mechanisms of the carp scale, revealing slight tensile anisotropy. Using in situ small-angle X-ray scattering during tensile testing, the toughening mechanisms of the scale, including the adaptive structural reorientation of lamellae as well as fibrillar sliding and elastic deformation, are quantified and compared with those of other fish scales.

INTRODUCTION
The integument of vertebrates is a continuous and heterogeneous layer covering the body surface that consists of two distinct strata: the epidermis and the dermis.1 The epidermis is mainly composed of keratin with a stratified structure, playing a crucial role as a permeability barrier and mechanical protection.1,2 Within it, many specialized cells are also able to develop non-skeletal epidermal scales, such as the ones in pangolin, and appendages such as feathers, hooves, hairs, claws, beaks, and nails, serving as multifunctional structural materials (armor, weapons, transportation tools, and thermal protection) to carry out assorted functions for organisms surviving in diverse environments.2,3 The deep counterpart to the epidermis is the dermis, which is collagen based.1 The ability to mineralize the dermis layer is a major innovation for many lower vertebrates and has enabled them to develop an efficient protective shelter, also called “dermal armor,” to prevent tissue damage caused by predatory attack.4,5 The fish scale is the most common one, providing effective protection for the overwhelming majority of fish species living from tropical to polar regions.

Based on the histological and histochemical characteristics, the scales of the current existing fish are mainly grouped into placoid, ganoid, elasmoid, dermal plates, and scutes.1 The placoid scale is characteristic of cartilaginous fish, such as sharks and rays. The ganoid scale, another primitive category, is also called the rhombic scale because of its diamond shape. There are mainly two types of ganoid scales existing currently: the quad-layered one, which is unique to Polypteriformes (bichirs and reedfish) and the double-layered one, which is commonly found on Lepisosteiformes...
The scutes and dermal plates are also rigid fish scales, being modified from ganoid scales. The scutes are composed of a thick bony base capped by a layer of non-collagenous hypermineralized tissue called hyaloine, while the dermal plates are composed exclusively of bone. The most common example of a fish possessing scutes is the armored catfish (Callichthyidae, Loricariidae, and Doradidae); dermal plates are also usually found in sturgeons, paddlefish, sticklebacks, boxfish, pufferfish, and seahorses. Due to their rigidity, scutes and dermal plates are only able to provide protection for limited fish species compared with the most evolved type, the elasmoid scale.

The elasmoid scale is, without question, the most prominent dermal armor among living vertebrates, including most of current fish species and some gymnophionan amphibians. The scales are thin collagenous plates embedded in the dermal pocket with roughly two-thirds of posterior surface overlapped by neighboring scales. Although considerably diverse in morphology and ornamentation (including both ctenoid and cycloid shapes), all elasmoid scales have a similar fundamental structure, which consists of three layers. In the developmental sequence, the first one to be formed is the external layer, which is an ornamented thin layer composed of minerals (mainly hydroxyapatites) and well-mineralized, randomly oriented collagen fibrils (woven fibrils). The second one to be formed is the tissue beneath the external layer, which is a thick basal plate composed of larger collagen fibrils forming a laminated plywood structure called elasmodine, from which the name of elasmoid scales originates. This plywood tissue (elasmodine) is slightly mineralized or completely unmineralized, resulting in great deformability of the scales. The last component to be formed is an extremely thin layer on top of the surface, which is composed of pure mineral and is referred to as the “limiting layer.” Based on previous studies on the structure and mechanical behavior of elasmoid scales in various fish species, including arapaima, red sea bream, bass, and tarpon, the protection mechanisms of elasmoid scales have been characterized: the well-mineralized outer layer, including the external layer and limiting layer, provides hardness against the bite from a predator’s attack, and the soft collagenous laminated base enables great deformability to dissipate the stored elastic energy through assorted mechanisms, such as lamellar rotation and separation, collagen fibril stretching/compression, fibrillar delamination, and bridging. The synergy of these mechanisms leads to outstanding scale toughness. By using primarily biopolymers and minerals in an intricate and ingenious hierarchical architecture, the elasmoid scales combine sufficient strength with excellent damage tolerance, properties that are often exclusive to synthetic materials, making these scales high-performance structural materials.

Here, we cast our interest on the elasmoid scales of the common carp (Cyprinus carpio). As a typical dermal armor, the scales protect carp and enable them to survive in various environments globally. Although the structure and some mechanical properties of carp scale have already been studied, the ultrafine structure of its collagenous lamellae has not been thoroughly investigated and the deformation mechanisms have not been quantified. In this work, we unravel the detailed structure of the carp scale, in particular using transmission electron microscopy (TEM) to characterize the ultrafine structure of the collagen fibrils in the lamellar inner core. We find that there are two collagenous frames comprising the plywood base of carp scale. The first one is composed of the collagenous lamellae made of isolated collagen fibrils; the arrangement of their orientations follows a helical pattern. The second one is a set of several layers composed of thinner collagen fibrils and forming a “sheet-like” structure oriented from the basal part toward the external surface.
layer. In addition to conventional tensile testing, we also employed in situ synchrotron small-angle X-ray scattering (SAXS) measurements during the tensile deformation to quantify the lamellar rotation and collagen fibril stretching/compression. The deformation mechanisms of this well-evolved elasmoid scale are identified and compared with the primitive elasmoid scales from the living fossil coelacanth. This evaluation of the principal features is expected to generate critical thinking regarding the design of new synthetic structural materials.

RESULTS AND DISCUSSION
Hierarchical Structure of the Scales of the Common Carp
The common carp (C. carpio) is a widespread freshwater teleost from the family Cyprinidae that originated from eutrophic waters in lakes and large rivers in Europe and Asia. The carp’s body is fully covered by imbricated/overlapped elasmoid scales anchored in the dermis pocket, as shown in Figure 1A. Each scale has an oval-like shape and two-thirds of its surface, the white part shown in Figure 1B, is covered by neighboring scales, which distribute stresses from a bite across a large volume of material and provide penetration resistance at a reduced weight.32 Figures 1C–1H manifest the surface characteristics of the scale’s exterior surface (named the limiting or osseous layer due to its high degree of mineralization) and the external layer (~40–50 μm thick). The well-mineralized limiting layer covers the external surface except at the periphery, which is different from the other type of elasmoid scale, such as the coelacanth fish scale, which contains pulp cavities in the external layer on the exposed region.33,34 The carp scale does not have these odontodes on the exposed (posterior) region. Instead, the posterior region (dark region shown in Figure 1B) has numerous mineral denticles scattered on it (Figure 1C). A close-up view of an isolated denticle, shown in Figure 1D, indicates that it is embedded in the dermis, which is composed of randomly oriented collagen fibrils (Figure 1E). The overlapped surface, presented in Figures 1F and 1G at different magnifications, shows two ornamentation characteristics: the circuli, as indicated by the dotted lines, and radii, indicated by the arrows. The circuli are concentric ridges or elevations on the superficial layer (dotted lines in Figures 1F and 1G); in some primitive fish species such as coelacanth or bowfin fish; they are in a radial position.8,33 The radii, as indicated by the arrows in Figures 1F and 1G, are the linear grooves starting from the focal region of the scale and intersecting circuli. These features aid the anchoring of the scales in the dermis pocket and improve their stability.35,36 A close-up view of the circuli (Figure 1G) shows the regularity of their arrangement, with an approximately equal spacing. Mineralized thin woven fibrils (~30–50 nm in diameter) are observed in the space between circuli ridges (Figure 1H), forming a randomly oriented loose network; this indicates that the limiting layer is extremely thin and discontinuous, leading to the partially exposed external fibrous layer. We compared the mineral contents of three types of elasmoid scales (carp, coelacanth, and arapaima) by performing thermogravimetric analysis (TGA); the results, provided in Figure S1, show that the carp scale has the lowest mineral content. The cross-sectional view (Figures 1I and 1J) also confirms our observations on the thickness of limiting and external layers; it also reveals the multilayered structure for carp scales, i.e., beneath the limiting layer and the external layer, which has been identified on the top surface, is the elasmodine (Bouligand) layer, which is the major component of the scale. The detailed structure of this layer will be unraveled later. Figure 1I also reveals that the elasmodine, which is the laminated inner structure of the scale, is composed of two groups of collagen layers with different mineralization.15 The outer (upper) portions of the collagenous lamellae are more mineralized than the lower portion, as further shown in Figure S2.
Figure 1. Surface Morphology and the Multilayered Structure of Carp Scales

(A) The common carp (Cyprinus carpio) from which the scales are extracted. Scale bar, 100 mm.

(B) The carp scale is a typical elasmoid scale with an oval shape; two-thirds of surface (the white part) are overlapped by neighboring scales. Scale bar, 10 mm.

(C–H) Scanning electron micrographs of the outer surface of the carp scale. (C) The mineralized denticles are scattered on the exposed part (dark region) and increase the surface roughness. Scale bar, 200 μm. (D) Most parts of the denticle are covered by the dermis and only the rounded tip is exposed. Scale bar, 50 μm. (E) The fibrous dermis tissue around the denticle tip. Scale bar, 1 μm. (F) The overlapped region of the carp scale suggests that the mineral ridges (circuli), indicated as dotted lines, are interrupted by cracks (radii), as indicated by arrows. Scale bar, 200 μm. (G) A close-up view reveals that the mineral ridges which define the circuli are regularly spaced. Scale bar, 20 μm. (H) The external layer, which is composed of woven mineralized fibrils (diameter ~30–50 nm), is exposed at the spaces between the ridges. Scale bar, 200 nm.

(I) Optical microscopy image of the cross-section of a carp scale. It clearly shows that the scales have a multilayered structure: the limiting layer, external layer, and elasmodine (Bouligand) layer. The outer elasmodine layer is more mineralized than the inner. Scale bar, 50 μm.

(J) Scanning electron micrograph of the region shown in (I). Scale bar, 50 μm.
Scanning electron microscopy (SEM) images of the oblique surface of freeze-fractured scales, provided in Figures 2A–2C, disclose the structural characteristics of the inner layer. The internal layer of the scale represents the stacking of collagenous lamellae to form a plywood-like arrangement that occupies a major volume fraction of the scale.

Arrangements of Lamella Orientation

As the arrows indicate in Figure 2A, the orientations of successive lamellae in the carp scale form a helicoidal pattern from interior to exterior, named the Bouligand structure. The rotation angle (determined by the SAXS experiments described below) is ~36°. A similar helical pattern is also observed in the basal plate of arapaima and tarpon scales. In contrast, the lamellae orientations in coelacanth scale form a double-Bouligand-type structure (shown in Figure 2D) in which one unit consists of two adjacent lamellae forming successive orthogonal plies (marked in different colors) in a helical pattern. Some other elasmoid scales such as those from red seabream, goldfish, and striped bass have an orthogonal arrangement for their collagenous plies, whereby two adjacent lamellae are always perpendicular.

Fibrillar Arrangement in Each Lamella

A close-up view of the assembly of collagen fibrils in one lamella, shown in Figures 2B and 2C, reveals that there is no grouping of fibril bundles formed in a carp scale.
Each lamella is composed of isolated collagen fibrils connected by sacrificial bonds (arrows in Figure 2C). A similar structure has been identified in most other elasmoid scales, such as arapaima, red seabream, tarpon, goldfish, and striped bass. This structure is different from the bundled collagen fibrils in coelacanth scales, shown bordered with dotted lines in Figure 2E, with the distinct characteristic d-spacing of collagen shown in Figure 2F.

Spacing between the Lamella Fibrils
Similar to other elasmoid scales including arapaima, tarpon, and striped bass, in the carp scale the isolated collagen fibrils in lamella are directly packed together and there is no noticeable substance between them, as shown in Figure 2B. In contrast, the space between the fiber bundles in the coelacanth scale is filled with a fibrous structure, as indicated in the inset of Figure 2E.

To further investigate the structural features of the collagenous lamellae of carp scales at the nanoscale, we characterized the internal layer by TEM; cross-sectional images are presented in Figures 3A–3D. We found that there is a secondary collagenous array perpendicular to the lamellae and aligned in the thickness direction, which has not been identified in previous studies on carp scales. Figure 3A shows three adjacent collagenous lamellae of the carp scale. Secondary collagenous arrays with bifurcations are embedded in it. Thin collagen fibrils thread through the thickness of lamellae; one of them is marked with a dotted line. Detailed views in Figures 3B–3D show that the threading collagen fibrils are relatively thin (d2 ≈ 20–40 nm) compared with those in the lamella (d1 ≈ 100 nm) shown in Figure 3D, and are interspersed within the lamella fibrils (Figures 3C and 3D); they do not group with other collagen fibrils, as observed in the coelacanth scale (Figures 3F and 3G). Such two-dimensional (2D) TEM images can be projected into a three-dimensional structure, in which the threading secondary collagen fibrils form interfaces that partition the collagen fibrils in the lamellae, as illustrated in the schematic in Figure 3E. More evidence showing that this vertical structure is not composed of single fibrils is provided in Figure S3. By comparison, the interbundle fibrils in the coelacanth scale are much more prevalent and robust (Figures 3F and 3G). Indeed, with the exception of fish from the family of Cyprinidae and Sarcopterygii, all other elasmoid scales lack such through-thickness collagen fibrils. Our hypothesis is that these threading through-thickness fibrils keep the lamellae together and prevent the scale from delaminating when subjected to external bending.

Tensile Behavior of Carp Scales
Based on our previous study of fish scales, the inner collagen core is under tension when the fish scale is under attack from the penetration of a predator’s teeth. Accordingly, to characterize their mechanical behavior, we performed uniaxial tensile tests on carp scales that were extracted from the mid-lateral region of a fresh carp. A representative tensile stress-strain curve is presented in Figure 4A; the insets indicate that the dog-bone-shaped specimens were cut along the transverse and longitudinal directions. The stress-strain curve shown in Figure 4 is from a transverse sample. Images of the specimen during testing at successive tensile strains of εt = 0, 0.06, 0.11, 0.20, 0.22, 0.25, 0.27, 0.34, and 0.43 are shown in Figures 4B–4J. In the linear region the scale deformed uniformly, as shown in Figures 4B–4D. Because of the greater elasticity of the collagen compared with the mineral constituents, the limiting and external layers cracked as tissue strain approached εt = 0.20, which was also accompanied by a small load drop on the stress-strain curve due to brittle cracking of the mineral in the external layer (as circled in Figure 4E). With continued stretching of the collagen fibrils, at εt = 0.22 the external layer started to separate.
Figure 3. The Secondary Fibrous Frame in the Carp and Coelacanth Scales

(A–D) Transmission electron micrographs of a carp scale with different magnifications. (A) The cross-section of three neighboring layers indicates there are a few very thin vertical fibrils going through the thickness direction. One of them is marked with a dotted line. Scale bar, 10 μm. (B) Higher-magnification image of the junction between two layers 1 and 2. Scale bar, 2 μm. (C) Close-up view of the cross-section reveals the delamination of the vertical fibril and that the collagen fibrils that constitute the lamella are tightly packed and squeezed into a polygonal shape. Scale bar, 500 nm. (D) High-magnification image shows that the thin vertical fibril is collagen, as indicated by the arrow. The diameter of the lamellar and vertical fibrils is denoted as $d_1$ and $d_2$, respectively. Scale bar, 100 nm.

(E) Schematic illustration of the structure of the vertical fibrils interspersed in the lamellae.

(F and G) Transmission electron micrographs of a coelacanth scale. The spaces between the fiber bundles (FBs) are filled with abundant interbundle fibrils (IBFs), which form a well-defined matrix to hold the Bouligand-type lamella. Scale bars, 5 μm (F) and 500 nm (G). The inset in (F) shows that the packing of the fibrils in the fiber bundle is not as tight as that in carp scales, and their cross-sections remain round in shape. Scale bar, 100 nm.
partially at the relatively weak interface from the stretched collagenous lamellae (Figure 4F), due to the modulus mismatch between the external mineralized layer and internal collagenous layers. At $\varepsilon_t = 0.25$, the external layer peeled off from the extreme tensile/deformed region of the collagen layer, as indicated by the dotted line in Figure 4G; this resulted in a more significant load drop. Beyond this point the engineering stress was almost zero, although the collagen fibrils still appeared to be stretching (Figures 4H–4J) under an ultralow sliding force, indicating that the lamellae were completely delaminated, consistent with material failure.

Figure 5 shows engineering stress-strain curves obtained from each of four conditions (seven tests per condition). The whole scale and samples without the limiting layer were tested in the longitudinal and transverse directions; the latter specimens...
were prepared by removing the highly mineralized limiting layer by mechanical polishing. The strength and elastic modulus in the longitudinal direction are slightly higher than in the transverse direction, although there is a much larger variation in the data for the longitudinal samples. Conversely, the total elongation tends to be smaller in the longitudinal direction. The resulting mechanical properties, specifically, the Young’s modulus, ultimate tensile stress, ultimate tensile strain and total energy dissipation, are listed in Table 1.

Such modest anisotropy in the tensile properties of carp scale was also observed in arapaima scales, although the property differences between orientations are more substantial. Striped bass scales also exhibit in-plane anisotropic tensile behavior but the trends are different, the strength in the transverse direction is higher but with a larger variation, which is opposite to the situation in the carp scales. In contrast, with its double-twisted Bouligand structure, the coelacanth scale is relatively isotropic in-plane. This is because the orientations of adjacent lamellae in the double Bouligand structure are perpendicular; the bilayer system forms in a twisting manner, resulting in more orientations of lamellae fibrils, which serves to increase isotropy. By comparison, the lamellae fibrils in the single-twisted Bouligand structure, e.g., in the carp scales, only have five orientations, which are equally spanned with 36° intervals, leading to a modest anisotropic tensile behavior.

Since the superficial and highly mineralized limiting layer in the carp scale is a small fraction of the overall thickness and the mineral content of the whole scale is low, there is almost no difference between the tensile properties of the whole scale and the sample without the limiting layer. In other elasmoid scales that possess a well-defined mineral layer, such as for the coelacanth, arapaima, and even
striped bass, the inner collagenous layer (without the highly mineralized limiting layer) has higher tensile strength than the whole scale. Akin to the arapaima scale, the thicker the limiting layer (also referred to as “mineral layer” in some sources), the more significant is the difference between the tensile properties for entire scale and the inner collagenous layer. Another factor that should not be ignored is that a minor portion of the external elasmodine (the Bouligand-type collagenous lamellae) was likely removed as well as the limiting layer. As indicated in Figure 1I, the elasmodine exhibits a gradient in mineralization with the outside being more mineralized, leading to different mechanical performance compared with the less mineralized innermost core. Since the removed external elasmodine may be inconsistent in different specimens, this could also contribute to the large variation in the tensile stress-strain curves for the samples without the limiting layer.

SEM images of the fracture surface after tensile failure of the carp scales, presented in Figure 6, show typical ductile tearing of the collagen. Since the external layer is vulnerable to peeling off from the collagen lamellae, there is no remaining external layer on the fracture surface (Figure 6F). Numerous straight fiber lamellae are revealed around the break region, which are due to the failure of the “sheet-like” secondary network enabling the collagenous lamella to separate. The fracture surface of relaxed broken fibers and fibrils after tensile failure (Figure 6E) indicates their stretching during extension of the bulk specimen. In the region adjacent to the fracture, other mechanisms can be identified, including fiber twisting, delamination, and bridging between fibrils; these are shown in Figures 6A, 6B, 6D, and 6H. The separation between the lamellae is shown in Figure 6C, together with their delamination due to the failure of the perpendicular sheets. Further evidence of interlamellar separation can be seen in Figure 6G in a region behind the fracture surface; groups of dotted lines in this image delineate the orientations of four adjacent lamellae. In general, these surfaces reveal that the lamellae fibrils are significantly stretched and delaminated when interlamellar separation occurs.

Summarizing, the toughening mechanisms act in concert to dissipate the stored energy over multiple length-scales and, as such, are responsible for the excellent damage tolerance of the carp scale. In comparison with previous studies, we have also confirmed that the total energy dissipation in the tensile testing of carp scale (5.83–6.80 MJ/m³) is definitively higher than that in other common collagen-based materials, e.g., human skin (2.44–4.28 MJ/m³), human cortical bone (1.8–2.5 MJ/m³), and arapaima fish scales (1.07–3.12 MJ/m³). Such excellent toughness is conferred by a sequence of multiple-scale toughening mechanisms including the separation between external and internal layer at the macroscale, the interlamellar separation and fiber splitting/twisting at the mesoscale, the processes of fibrillar

| Table 1. Uniaxial Tensile Test Results in both Longitudinal and Transverse Directions for the Whole Carp Scale and the Collagen Layer Only |
|-----------------|-----------------|-----------------|-----------------|-----------------|
|                 | Whole Scale     | Collagen Layer  |
|                 | Longitudinal    | Transverse      | Longitudinal    | Transverse      |
| Young’s modulus (MPa) | 521.5 ± 125     | 463.5 ± 92.3   | 526.0 ± 126.1   | 406.4 ± 95.5   |
| Ultimate tensile stress (MPa) | 57.4 ± 14.5     | 41.4 ± 12.8   | 57.3 ± 17.8    | 51.7 ± 7.00   |
| Ultimate tensile strain (mm/ mm) | 0.193 ± 0.023     | 0.206 ± 0.047 | 0.196 ± 0.043   | 0.201 ± 0.005 |
| Energy dissipation (MJ/m³)    | 6.79 ± 1.75     | 5.83 ± 2.90   | 6.8 ± 2.31     | 6.22 ± 2.23   |
delamination and bridging at the microscale, and overall the adaptive structural reorientation of the collagen in response to external loading.

**In Situ Synchrotron SAXS Analysis of Toughening Mechanisms in Carp Scales**

Toughening by adaptive structural reorientation, whereby the majority of the fibers or fibrils rotate in order to carry more of the load, is not uncommon in collagenous materials; notable examples are arapaima fish scales\(^{16}\) and the dermis of mammals.\(^{41}\) Recent studies\(^{16,35}\) on the mechanical behavior of collagenous tissue show that *in situ* SAXS is a powerful tool to precisely characterize the deformation and reorientation of collagenous lamellae under load.\(^{42}\) This is attributed to the unique structure of collagen fibrils. In the collagen-based inner core of elasmoid scale, each collagen fibril is composed of a quarter-staggered array of molecules with a certain amount of mineral nanoparticles filling the gap regions situated between the heads and tails of the collagen molecules,
forming alternating bands (i.e., the distance between the center of two gaps); this is known as the $d$-spacing or $d$-period. When X-rays are projected onto the tissue, the collagen fibrils act as a molecular diffraction grating due to the periodicity of the collagen-mineral composite structure, generating a well-defined diffraction pattern, which can be used for quantitative analysis. Figure 7A shows the experimental setup that we used. Samples in both longitudinal and transverse directions of the scales were examined by synchrotron X-rays during tension testing with a series of 2D SAXS patterns obtained, as shown in Figures 7C and 7D. The dark blue rectangular mesh in the SAXS pattern is due to the space between the modules of the Pilatus detector. Representative SAXS patterns generated by the unloaded longitudinal samples are shown in Figure 7C. The interaction of the X-rays with collagenous lamellae creates several sets of concentric arcs in the diffraction pattern. Based on Bragg’s law, the distance between the beam center and the first-order arc is the reciprocal of the $d$-spacing of collagen fibrils (Figure 7B), which is $63.2 \pm 1.1 \text{nm}$ in our carp scale specimens. A detailed explanation of this reciprocal relationship is provided in Figure S6. Based on the change in this distance in each diffraction ring, the elastic strain in the collagen fibrils can be quantified. The azimuthal angle of the arcs, referred to as $\Psi$, indicates one distinct orientation of the collagen fibrils, as illustrated in Figure 7B. The deformation of the collagen fibrils under uniaxial tension loading ($\Psi = 0$) can be quantitatively analyzed for the sequence of the SAXS pattern. Further details about the analysis are provided below.

Figure 7D shows a typical tensile stress-strain curve with five SAXS patterns revealing the deformation of the collagen fibrils under loading at tissue strains $\varepsilon_t$ of 0, 0.06, 0.10, 0.14, and 0.18. A close-up view of the first-order Debye-Scherrer ring coupled with the real-time pictures at the initial state ($\varepsilon_t = 0$) and fully stressed state ($\varepsilon_t = 0.18$) are provided in the insets. When the scale is fully relaxed, the diffraction pattern is composed of round arcs with some preferred orientations. As the applied tensile load increases, the diffraction pattern evolves accordingly. When the tensile stress reaches its maximum at $\varepsilon_t = 0.18$, the specimen is fully stretched and the diffraction pattern changes from a round shape (before tension) to a hexagon-like shape at the maximum stress, indicating significant anisotropic strains in the structural evolution of the inner lamellar layer.

Reorientation of the Collagen Fibrils under Load
To quantify the real-time reorientation of the collagen fibrils, we generated plots of intensity versus azimuthal angle (orientation of the collagen fibrils) at five strains (0, 0.06, 0.10, 0.14, and 0.18) from the 2D SAXS pattern during tension testing, as shown in Figure 8A. The data on the fractured sample at a strain of 0.22 are also included. Prior to testing (at $\varepsilon_t = 0$), the single Bouligand structure of the carp scale showed highly oriented collagen fibrils with an included angle of $\sim36^\circ$ as the high-intensity peaks ($-78^\circ$, $-46^\circ$, $-9^\circ$, $27^\circ$, $61^\circ$) indicated in the plot ($\varepsilon_t = 0$) in Figure 8A. This unique twisting angle is very consistent in this single-twisted Bouligand structure, which was confirmed by investigation of several specimens (Figure S4). As the applied load is increased, the collagen fibrils at large angles (i.e., $\Psi = -78^\circ$, $61^\circ$) reorient further away from the loading direction with increasing intensity (shown by arrows); several peaks (dotted circles) appearing at the angles of $-65^\circ$ to $-55^\circ$ reveal that the collagen lamellae reorient toward, and delaminate along, that orientation. As the loading is increased, lamellae at angles closer to the loading direction ($\Psi = -9^\circ$ to $\Psi = 27^\circ$) reorient and distribute over a broad range (between $-40^\circ$ and $40^\circ$) toward the loading direction (shown by the stars and triangles in Figure 8A). After the scale is broken (at $\varepsilon_t = 0.22$), the peaks reappear but at more orientations than in the original state, indicating that a significant portion of the lamellae have reoriented and been delaminated, i.e., the scale has experienced significant permanent...
Figure 7. In Situ SAXS Patterns of Carp Scales Taken in Real Time during a Uniaxial Tensile Test

(A) Experimental setup. Tensile specimens are extracted from the overlapped region of the scales along both longitudinal and transverse directions, and the mineral layer is removed mechanically. The prepared specimen is exposed to the synchrotron X-rays in the beamline at the Advanced Light Source during in situ uniaxial tensile tests, with the scattering patterns recorded as a function of the tissue strain in real time.

(B) Schematic showing the reciprocal relationship between the radius of the first-order arc (q) and the d-spacing of the collagen fibrils, based on Bragg’s law.

(C) Pattern of a fully relaxed longitudinal sample shows several sets of concentric arcs, each of them caused by the diffraction of collagen fibrils well aligned in certain directions, which are the lamellar orientations in the fish scale. The loading direction is defined as the azimuthal angle (ψ) equal to 0°.

(D) In situ SAXS patterns at five different tissue strains (εt = 0, 0.06, 0.10, 0.14, and 0.18) in the tensile test of a longitudinal sample. The insets in the patterns of the fully relaxed state (εt = 0) and highest stress state (εt = 0.18) are the close-up views of first several Debye-Scherrer rings with the real-time images of the tissue. The diffraction pattern changes from a round circle with six equal sectors to a hexagon-like shape, revealing the inhomogeneous deformation of the fibrils in different orientations.
deformation and the collagenous lamellae are no longer oriented with a specific included angle of $36^\circ$.

Compared with the single Bouligand structure of carp scales, the toughening mechanisms in the double Bouligand structure of the coelacanth scale are simpler, as shown for comparison in Figure 8B. Since the structure is formed by orthogonal bilayers twisted in a Bouligand pattern, collagen lamellae are aligned in more orientations with no clear high-intensity peaks. As the load increases, collagen fibrils ranging between $\psi = -60^\circ$ and $\psi = 40^\circ$ gradually rotate toward the loading direction (shown by arrows at $0^\circ$) due to interfibrillar sliding in response to the shear stresses; the ones at large angles ($-90^\circ < \psi < -60^\circ$, $40^\circ < \psi < 90^\circ$) rotate away from the loading direction and delaminate/separate significantly (shown by a lower intensity). By rotating toward the loading direction, the projected length of the fibrils increases. This rotation has an analogy in the plastic deformation by single slip of metallic monocrystals. This mechanism was explained and quantified by Yang et al.\textsuperscript{17} for the arapaima scales and modeled by molecular dynamics. On the other hand, the rotation away from the direction of loading, for fibrils oriented at large angles to it, can be explained by separation of the fibrils in the lamellae produced by tension. This creates curved fibrils with varying angles to the tensile axis.

**Quantification of the Deformation of Lamellae of Collagen Fibrils**

For the determination of the fibrillar strains, the diffraction circles were evenly divided into 17 sectors with a span of $10^\circ$ from $\psi = -90^\circ$ to $\psi = 70^\circ$, and the average radius $r$ of each sector was calculated. The radius $r$ is inversely proportional to the characteristic $d$-spacing through Bragg’s law, and thus the $d$-spacing of the collagen fibrils in each sector was obtained. The change in $d$-spacing in each sector was then divided by the initial, unstressed $d$-spacing of $d_0$ to define the collagen fibril strain $\varepsilon^f_i$. The strains are plotted as a function of the azimuthal angle in Figure 9A. Calculations were performed for the five externally imparted strain levels (tissue strains)
during tensile deformation: 0, 0.06, 0.10, 0.14, 0.18, and 0.22 (at failure). For comparison, a similar plot for the coelacanth scale is presented in Figure 9B.  

Figure 9A shows that the imparted tissue strain results in tensile fibrillar strains for small angles ($\Psi < 40^\circ$). This is most pronounced for $\Psi = 0^\circ$ and $\Psi = 10^\circ$ and increases with tissue strain. However, the fibrillar strain is significantly lower (by approximately one-half) than the tissue strain by virtue of interfibrillar sliding. The maximum tensile fibril strain, when the tissue is fully stretched, is about 0.1; this is considerably lower than the global tissue strain, 0.18, indicating that significant fibrillar sliding has taken place. The same mechanism was also identified in both coelacanth and arapaima scales; however, the fibril extensibility of carp scale ($\varepsilon_f \sim 0.08$) is higher than that of the arapaima scale ($\varepsilon_f \sim 0.05$), but lower than the coelacanth scale ($\varepsilon_f \sim 0.12$). A similar effect also exists in rabbit dermis. Such sliding between collagen fibrils in these collagenous tissues, e.g., fish scales and mammal dermis, is permanent.

For the lamellae oriented at large angles ($-90^\circ < \Psi < -40^\circ$ and $40^\circ < \Psi < 90^\circ$), shown in Figure 9A, the fibrils are compressed. This reduction in length is associated with the tensile stretching of the lamellae at small angles $\Psi$ with the tensile direction. These lamellae exhibit a reduction in their lateral dimensions due to a Poisson’s ratio effect. This reduction in the lateral dimension causes a compression in the lamellae at large angles, which are connected to them. Thus, negative strains are apparent, shown in Figure 9A, for $\Psi$ between $\pm 40^\circ$ and $\pm 90^\circ$. This unique deformation
Figure 10. Schematic Rendition Showing the Mechanisms for the Adaptive Deformation of the Carp Scale under Tensile Loading

The purple arrows indicate the loading direction.
(A) The mineral layer cracks and separates from collagenous base.
(B) The lamellae along the tensile direction are stretched ($d > d_0$) and finally break as loading increases.
(C) Interfibrillar sliding causes the lamellae to rotate toward the loading direction (as the black arrow indicates) and the collagen fibrils are also stretched, with increased $d$-spacing ($d > d_0$).
(D) The lamellae initially oriented far from tensile axis are compressed ($d < d_0$) and rotate away (as the black arrow indicates). With further increase in loading, the lamellae fail by interfibrillar delamination.
(E) The original collagenous lamellae form a highly ordered Bouligand-type structure with a rotation angle of 36°.
(F) The overall deformed collagenous lamellae with applied tensile load. Combining all the adaptive mechanisms in (B) to (D) and the separation and shear between lamellae, the fibrous base confers the scale with significant deformability to adapt to the applied load.
mechanism for carp scales is also confirmed by measurements made on several specimens, along different directions, as illustrated in Figure S5.

These SAXS results and their analysis, presented in Figures 8 and 9, can be summarized in the schematic drawing in Figure 10, which shows the sequence of mechanisms experienced by the collagenous lamellae in the carp scale based on both lamellar reorientation (quantified in Figure 8A) and deformation (Figure 9A). The first stage during tensile loading is the cracking and peeling off of the highly mineralized external layer (Figure 10A), which is confirmed by our observations in Figures 4E–4G. The lamellae with orientations close to the loading axis (−40° < Ψ < 40°) are subjected to tensile strains, resulting in an increase in length with increasing d-spacing (d > d₀, Figures 10B and 10C) and reorient toward the loading direction by interfibrillar shear (Figure 10C). These mechanisms, which are proposed on the basis of our analysis of the SAXS experiments (from Figures 8A and 9A), are confirmed by our observations on the fracture surface of the tensile specimens presented in Figure 6, which specifically show fibrillar shearing and delamination (Figure 6A) and relaxation of stretched fibers after tensile failure (Figure 6E). There is some difference in the fibrillar strain distribution between the carp and coelacanth scale, shown in Figures 9A and 9B. Since the threading fibrils in carp scale are not as profuse as the ones in coelacanth scale, the lamellae are more vulnerable to failure by delamination as illustrated in Figure 10D. This delamination is characterized by SEM in Figures 6B and 6C. Interestingly, a transitional range exists where the fibril deformation changes from tension to compression, i.e., the fibrils initially oriented from −40° to −30° and from 40° to 50°. This indicates that deformation takes place exclusively by shear on some fibrils within these two ranges. The lamellae initially oriented far from tensile axis are compressed (d < d₀, Figure 10D) and rotate away (as the black arrow indicates in Figure 10D) because of the Poisson’s effect. The fibrils strained at εₜ = 0.22, when the tissue completely breaks, do not recover their original value (εₜ = 0), indicating the retention of tensile and compressive plastic deformation when the scale is stretched to fracture. Such plastic deformation is also evident on the ductile fracture surfaces shown in Figures 4J and 6A.

By combining this sequence of deformation mechanisms discussed above, the Bouligand-type lamellar inner base of the carp scale adapts to the applied load by lamellar rotation, fibrillar deformation and sliding, and interlayer separation, as illustrated in Figure 10E and 10F, and as such is capable of dissipating stored elastic energy effectively and conferring to the carp scale excellent toughness.

Conclusions
We have characterized the detailed structure of the modern elasmoid scales of the common carp and determined their mechanical properties. Using synchrotron SAXS, the lamellar rotation and fibril deformation have been quantified to identify the salient toughening mechanisms. To improve our understanding of the prevailing mechanisms, we have further compared our results for the carp scale with those of the primitive elasmoid scales from the coelacanth. Based on this work, several conclusions on the prime structural components of this damage-tolerant natural material can be made.

- With respect to the structure of carp scales, the external layer is composed of highly mineralized woven fibrils covered by an extremely thin and discontinuous mineral layer. The inner core is a highly ordered “plywood-like” structure, consisting of superimposed collagenous lamellae, successively rotated by 36° to form a typical single-twisted Bouligand structure. The small amount of
secondary “threading” fibrils along the thickness direction is presumed to ensure the integrity of the scales.

- The tensile behavior of the carp scale, which represents the behavior of the inner collagenous layer under penetration (i.e., from attack by a predator), was found to have higher strength and toughness (by ~20%) in the longitudinal, as compared with the transverse, direction. A synergy of toughening mechanisms including external layer separation, interlamellar separation, fiber twisting and splitting, fibrillar delamination and bridging, and fibrillar reorientation, all act in concert to improve the damage tolerance of the scales.

- In situ SAXS measurements during mechanical straining of the carp scales demonstrate adaptive structural reorientation, in that the lamellae oriented closer to the tensile direction rotate toward it to carry more load while stretched and the ones oriented far from the loading direction rotate away from it while compressed. However, the collagen fibrils present two local maxima strains at azimuthal angles of $\Psi = -20^\circ$ and $\Psi = 10^\circ$ to the tensile axis. Compared with the single Bouligand structure of carp, the double Bouligand structure containing the secondary “threading” bundles of coelacanth scales exhibits higher integrity so that the collagen fibrils reach to one clear maximum fibrillar tension strain under load.

- These deformation mechanisms operate together to render the carp scale an excellent dermal armor, and may well provide further inspiration for the design of advanced synthetic structural materials with unprecedented toughness and penetration resistance.

**EXPERIMENTAL PROCEDURES**

**Resource Availability**

**Lead Contact**

Further information and requests for resources and reagents should be directed to and will be fulfilled by the Lead Contact, Robert O. Ritchie (roritchie@lbl.gov)

**Materials Availability**

The scale samples were collected from a common carp (C. carpio) with a length of 670 mm, which was purchased from Ranch 99 Market (San Diego, CA, USA). The scales were extracted from their pockets with tweezers and preserved in deionized water. They varied from 20 to 40 mm in length but the ones that we used were all ~30 mm in length and extracted from the mid-lateral region. For comparison, coelacanth scales were obtained from a *Latimeria chalumnae* with a length of 950 mm in the Marine Vertebrate Collection of Scripps Institution of Oceanography, University of California, San Diego; the specimen was preserved in 80% isopropanol since its collection from the Grand Comore Island in 1973. The scales were peeled off from the linea lateralis below the anterior dorsal fin.

**Data and Code Availability**

Data supporting the findings of this paper are available from the corresponding authors upon reasonable request.

**Structural Characterization**

The structure of the scales and the fracture surfaces after tensile testing were all characterized in an FEI Quanta 250 and FEI Apero scanning electron microscope (FEI, Hillsboro, OR). To obtain oblique fracture surfaces, we immersed scales in liquid nitrogen for 30 s and immediately fractured them using forceps. The samples were first immersed in the 2.5% glutaraldehyde for 1 h to fix the structure and then dehydrated with an ascending series of ethanol (30, 50, 70, 90, 95
and 100 vol % twice), before being dried using an automatic critical point dryer (AutoSamdri 815A; Tousimis, Rockville, MD). All dried samples were then sputter coated with iridium using an Emitech K575X sputter coater (Quorum Technologies) before observation. TEM images were taken by a Technai 12 (Spirit, 120-kV) transmission electron microscope (FEI).

Preparation of TEM Specimens
The carp and coelacanth scales were first cut into small strips, with a length of 5 mm and a width of 2 mm, then immersed into 2.5% glutaraldehyde in 0.15 M sodium cacodylate buffer (pH 7.4) for 1 h to fix the structure. The fixed specimens were stained with 1% OsO4 solution with 8% potassium ferrocyanide in 0.5 M sodium cacodylate buffer overnight at room temperature. The scales were then stained with 2% aqueous uranyl acetate for 12 h and subsequently dehydrated with an ascending ethanol series (50%, 70%, 90%, and 100% twice), followed by a 1:1 ratio of 100% ethanol and 100% acetone, and finally 100% acetone. The fully dehydrated specimens were embedded in Spurr’s low-viscosity resin and polymerized at 60 °C for 72 h. The obtained blocks were then sectioned parallel to the vertical cross-section before ultrathin slices with thickness of ~80 nm were generated using a Leica Ultracut UCT ultramicrotome (Leica) and a Diatome diamond knife (Diatome). Ultramicrotomed sections were then placed on copper grids for TEM observation and post stained with Sato lead for 1 min before final examination.

Uniaxial Tensile Testing
Dog-bone-shaped tensile specimens of carp scales with dimensions of 20 × 3 mm were cut along longitudinal and transverse directions from carp scales with a thickness of between 0.1 and 0.2 mm. The outer layers of some samples were removed using silicon carbide polishing paper of 180# to 2500#, leading to a final thickness less than 0.1 mm. To prevent sliding during tensile testing, we glued the ends of the samples between sandpaper sheets using cyanoacrylate glue, resulting in a gauge length of ~10 mm. Uniaxial tensile tests were performed on an Instron 3342 mechanical testing machine (Instron, Norwich, MA) with a load cell of 500 N at a strain rate of 10^{-3} s^{-1}; samples were tested immediately after being removed from deionized water, where they had been kept prior to testing. Values of Young’s modulus were estimated by linear fitting the stress-strain curve at the early stage of the tensile testing (in the linear portion). We also determined the ultimate tensile stress, uniform elongation (strain at maximum stress), and the toughness, the latter being measured in terms of the area beneath the uniaxial stress-strain curves.

In Situ SAXS during Mechanical Testing
Using identical specimen preparation procedures as for the uniaxial tensile tests, specimens for in situ SAXS measurements were cut in both longitudinal and transverse directions from complete carp scales after the mineral layer had been removed. The resulting specimens were loaded in tension while simultaneously exposing them to synchrotron X-rays at beamline 7.3.3 at the Advanced Light Source synchrotron radiation facility (Lawrence Berkeley National Laboratory, Berkeley, CA). The tensile tests were performed using a Linkam TST-350 tensile stage (Linkam Scientific Instruments) with a 200-N capacity load cell to measure the force. Such an experimental setup allows for SAXS data collection to be recorded in real time with the measurement of the load-displacement curve. The mechanical tests were performed at room temperature and a strain rate of 1.0 × 10^{-3} s^{-1}.

During mechanical testing, a high flux of synchrotron X-rays was focused on the tensile specimen and scattered by the periodicity of the collagenous lamellae, which in
essence serves as a macroscopic strain gauge; this generates a diffraction pattern composed of several concentric arcs on the screen, shown in Figure 7C. In each set of concentric arcs, the arc closest to the beam center (the first-order arc) represents the distance $q$ from the beam center, which has a reciprocal relationship with the $d$-spacing of the collagen fibrils, whereas the rest of the concentric arcs are higher-order arcs, corresponding to the harmonics of this $d$-period. The azimuthal angle of each set of arcs, referred to as $\Psi$, indicates one distinct orientation of the collagen fibrils in the scale, as illustrated in the inset of Figure 7B. Uniaxial tensile loading was applied along the $\Psi = 0^\circ$ direction. Therefore, the change of in the value of $\Psi$ $q$ and $q$ $\Psi$ for all sets of the first-order arcs at different tissue strains can quantify, respectively, the rotation of the collagenous lamellae and the deformation of strain in the collagen fibrils in real time during uniaxial tensile extension.

A Pilatus3 2M detector (Dectris, Baden, Switzerland) was used to collect the SAXS data. The detector was located at the largest allowable distance, around 4 m, to detect fine changes in the position of the collagen peaks. The sample was exposed to X-rays for 0.5 s with intervals of $\sim$5 s during the mechanical test; this radiation dosage was sufficiently low so as not to affect the structure and properties of the scales.

Quantification of Fibril Orientation Using SAXS

On examining the SAXS spectra patterns, 360 even sectors were made on the Debye-Scherrer rings with a span of $1^\circ$, starting from $\Psi = 0^\circ$ to $\Psi = 360^\circ$; as such, a sector graph was generated by the polar transformation of the 2D diffraction pattern using the software IGOR Pro (Wavemetrics) in conjunction with the custom macro NIKA (Jan Ilavsky, Argonne National Laboratory, IL, USA). Based on the obtained square map of intensity versus pixel, the 1D graphs, specifically the plots of the integrated intensity as a function of azimuthal angle, were created by using the “Image line profile” tool in NIKA.

Quantification of Fibril Strains

The strains in the collagen fibrils were measured from the X-ray data in terms of the change in 1D peak position determined from the plot of integrated diffraction intensity as a function of $d$-spacing. By using the same software as in the quantification of fibril orientation, the sample detector distance and beam center were calibrated with the 2D diffraction pattern of a silver behenate standard. To convert the 2D SAXS data into 1D peaks, we evenly made 17 sectors on the upper half of the Debye-Scherrer rings, starting from $\Psi = -5^\circ$ to $\Psi = 165^\circ$, with a span of $10^\circ$ in each sector; the integrated intensity over the diffraction arc in each sector was then radially averaged to obtain the relationship between the intensity peaks and the radial distance of the arc, $q$. Based on the numerical relationship between $q$ and the $d$-spacing, plots of integrated intensity as a function of the $d$-spacing for all 17 sectors were generated by the software; intensity peaks were fit to an exponential Gaussian function and a linear background to precisely locate the peak positions. The strain in the collagen fibrils was measured as the change in position of the center of the first-order collagen peak, normalized by the unstressed state.

SUPPLEMENTAL INFORMATION

Supplemental Information can be found online at https://doi.org/10.1016/j.matt.2020.05.011.

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AUTHOR CONTRIBUTIONS

DECLARATION OF INTERESTS
The authors declare no competing interests.

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